

derivation of Ewald Summation

$$z = (x, y, z) = (r, \theta, \varphi)$$

z点受到的所有电荷产生的电势:

$$\Phi(z) = \frac{1}{4\pi\epsilon_0} \sum_n \sum_{j=1}^N \frac{q_j}{|z - z_j + nL|}$$

z点受到的除i之外的电荷产生的电势:

$$\Phi_{[i]}(z) = \frac{1}{4\pi\epsilon_0} \sum_n \sum_{j=1(\neq i)}^N \frac{q_j}{|z - z_j + nL|}$$

$$\Phi_{[i]}(z) = \frac{1}{4\pi\epsilon_0} \sum_n \sum_{j=1(\neq i)}^N \int \frac{\rho_j(z')}{|z - z' + nL|} d^3z'$$

电荷j产生的电荷密度:

$$\begin{aligned} \rho_j(z) &= q_j \delta(z - z_j) \\ &= q_j \left(\delta(z - z_j) - G(z - z_j) + G(z - z_j) \right) \\ &= q_j \underbrace{\left(\delta(z - z_j) - G(z - z_j) \right)}_{\rho_j^s(z)} + q_j \underbrace{G(z - z_j)}_{\rho_j^l(z)} \end{aligned}$$

高斯分布:

$$G(z) = \frac{\alpha^3}{\pi^2} e^{-\alpha^2 r^2}$$

z点受到的电荷密度 ρ 产生的电势方程 (泊松方程):

$$\nabla^2 \Phi(z) = -\frac{\rho(z)}{\epsilon_0}$$

$$\nabla^2 \Phi_j^l(z) = -\frac{\rho_j^l(z)}{\epsilon_0}$$

$$\frac{1}{r} \frac{\partial^2}{\partial r^2} (r\Phi) = -q_j \frac{\alpha^3}{\epsilon_0 \pi^2} e^{-\alpha^2 r^2}$$

$$\frac{\partial^2}{\partial r^2} (r\Phi) = -q_j \frac{\alpha^3}{\epsilon_0 \pi^2} r e^{-\alpha^2 r^2}$$

$$\frac{\partial}{\partial r} (r\Phi) = q_j \frac{\alpha}{2\epsilon_0 \pi^2} e^{-\alpha^2 r^2} + C_1$$

$$r\Phi = q_j \frac{\alpha}{2\epsilon_0 \pi^2} \int_0^r e^{-\alpha^2 r^2} dr + C_1 r + C_2$$

$$\Phi = q_j \frac{\alpha}{2\epsilon_0 \pi^2 r} \int_0^r e^{-\alpha^2 r^2} dr + C_1 + C_2 \frac{1}{r}$$

$$= q_j \frac{\alpha}{2\epsilon_0 \pi^2 r} \frac{1}{\alpha} \int_0^r e^{-\alpha^2 r^2} d(\alpha r) + C_1 + C_2 \frac{1}{r}$$

$$= q_j \frac{1}{4\pi\epsilon_0 r} \text{erf}(\alpha r) + C_1 + C_2 \frac{1}{r}$$

$$\lim_{r \rightarrow \infty} \text{erf}(r) = 1$$

$$\lim_{r \rightarrow \infty} \Phi(z) = q_j \frac{1}{4\pi\epsilon_0 r} \text{ (Coulomb's law)}$$

$$\Rightarrow C_1 = C_2 = 0$$

电荷j产生的电势:

$$\therefore \Phi_j^l(z) = q_j \frac{1}{4\pi\epsilon_0 r} \text{erf}(\alpha r)$$

$$\Phi_j^s(z) = q_j \frac{1}{4\pi\epsilon_0 r} \text{erfc}(\alpha r)$$

z点受到的除i之外的电荷产生的电势:

$$\Phi_{[i]}(z) = \Phi_{[i]}^s(z) + \Phi_{[i]}^l(z)$$

$$= \sum_n \sum_{j=1(\neq i)}^N \Phi_j^s(z) + \sum_n \sum_{j=1(\neq i)}^N \Phi_j^l(z)$$

$$= \sum_n \sum_{j=1(\neq i)}^N q_j \frac{\text{erfc}(\alpha|z - z_j + nL|)}{4\pi\epsilon_0|z - z_j + nL|}$$

$$+ \sum_n \sum_{j=1(\neq i)}^N q_j \frac{\text{erf}(\alpha|z - z_j + nL|)}{4\pi\epsilon_0|z - z_j + nL|}$$

real part(cutoff内直接计算):

$$\Phi_{[i]}^s(z) = \sum_n \sum_{j=1(\neq i)}^N q_j \frac{\text{erfc}(\alpha|z - z_j + nL|)}{4\pi\epsilon_0|z - z_j + nL|}$$

wave part(倒易空间计算):

$$\Phi_{[i]}^l(z) = \sum_n \sum_{j=1}^N q_j \frac{\text{erf}(\alpha|z - z_j + nL|)}{4\pi\epsilon_0|z - z_j + nL|}$$

self correction:

由于wave part的FT计算时无法单独忽略自身(注意wave part公式里没有“ $\neq i$ ”), 自己和自己的作用也会被包含进去, 因此最后需要减掉这一项

$$\begin{aligned} \Phi_i^{self} &= \lim_{z \rightarrow 0} q_j \frac{\text{erf}(\alpha z)}{4\pi\epsilon_0 z} \\ &= \frac{q_j}{4\pi\epsilon_0} \lim_{z \rightarrow 0} \frac{\text{erf}(\alpha z)}{z} \\ &= \frac{q_j}{4\pi\epsilon_0} \lim_{y \rightarrow 0} \frac{2\alpha}{\sqrt{\pi}} e^{-y^2} \\ &= \frac{\alpha q_j}{2\pi^2 \epsilon_0} \end{aligned}$$

$$\begin{aligned} y &= \alpha z \\ z &= \frac{1}{\alpha} y \\ dz &= \frac{1}{\alpha} dy \\ \frac{d}{dz} \text{erf}(\alpha z) &= \alpha \frac{d}{dy} \text{erf}(y) = \frac{2\alpha}{\sqrt{\pi}} e^{-y^2} \end{aligned}$$

wave part

$$\vec{z} = (x, y, z)$$

电荷j产生的施加给z点的电荷密度:

$$\rho_j^l(\vec{z}) = q_j G(z - z_j)$$

z点的所有电荷施加的长程电荷密度:

$$\rho^l(\vec{z}) = \sum_n \sum_{j=1}^N q_j G(\vec{z} - \vec{z}_j + nL)$$

高斯分布:

$$G(\vec{z}) = \frac{\alpha^3}{\pi^2} e^{-\alpha^2 |\vec{z}|^2}$$

泊松方程:

$$\nabla^2 \phi(z) = -\frac{\rho(z)}{\epsilon_0} \longrightarrow |\vec{w}|^2 \tilde{\phi}^l(\vec{w}) = \frac{\tilde{\rho}^l(\vec{w})}{\epsilon_0} \quad (\text{证明见附录[1]})$$

$$\tilde{\phi}^l(\vec{w}) = \frac{e^{-\frac{|\vec{w}|^2}{4\alpha^2}}}{|\vec{w}|^2 \epsilon_0} \sum_{j=1}^N q_j e^{-i\vec{w}\vec{z}_j}$$

$$\phi^l(\vec{z}) = \frac{1}{L^3} \sum_{\vec{w} \neq 0} \tilde{\phi}(\vec{w}) e^{i\vec{w}\vec{z}}$$

$$\begin{aligned} \tilde{\rho}^l(\vec{w}) &= \int_V \rho^l(\vec{z}) e^{-i\vec{w}\vec{z}} d^3\vec{z} \\ &= \int_V \sum_n \sum_{j=1}^N q_j G(\vec{z} - \vec{z}_j + nL) e^{-i\vec{w}\vec{z}} d^3\vec{z} \\ &= \sum_{j=1}^N q_j \int_{R^3} G(\vec{z} - \vec{z}_j) e^{-i\vec{w}\vec{z}} d^3\vec{z} \\ &= \frac{\alpha^3}{\pi^2} \sum_{j=1}^N q_j \int_{R^3} e^{-\alpha^2 |\vec{z} - \vec{z}_j|^2} e^{-i\vec{w}\vec{z}} d^3\vec{z} \\ &= e^{-\frac{|\vec{w}|^2}{4\alpha^2}} \sum_{j=1}^N q_j e^{-i\vec{w}\vec{z}_j} \quad (\text{这一步推导过程见附录[4]}) \end{aligned}$$

exafmm/ewald.h

potential def:

$$\phi(z) = \sum_n \sum_{j=1}^N \frac{q_j}{|z - z_j + nL|}$$

real part:

$$\phi_{[i]}^s(z) = \sum_n \sum_{j=1(\neq i)}^N q_j \frac{\text{erfc}(\alpha|z - z_j + nL|)}{|z - z_j + nL|}$$

$$F_{[i]}^s(z) = \nabla \phi_{[i]}^s \quad \begin{array}{l} \text{结果及证明在附录[2]} \\ \text{(实际上这个负电场, 力还需要乘-q)} \end{array}$$

dft:

$$\tilde{\rho}^l(\vec{w}) = \sum_{j=1}^N q_j e^{-i\vec{w}\vec{z}_j}$$

wave part:

$$\tilde{\phi}^l(\vec{w}) = \frac{4\pi e^{-\frac{|\vec{w}|^2}{4\alpha^2}}}{L^3 |\vec{w}|^2} \tilde{\rho}^l(\vec{w})$$

idft:

$$\phi^l(\vec{z}) = \sum_{\vec{w} \neq 0} \tilde{\phi}(\vec{w}) e^{i\vec{w}\vec{z}}$$

$$F^l(\vec{z}) = \nabla \phi^l(\vec{z}) \quad \begin{array}{l} \text{结果及证明在附录[3]} \\ \text{(实际上这个负电场, 力还需要乘-q)} \end{array}$$

self correction:

$$\phi_i^{self} = \frac{2\alpha q_j}{\sqrt{\pi}}$$

$$F_i^{self} = \nabla \phi_i^{self} = 0$$

附录[1]: reciprocal空间的泊松方程

$$f = e^{i\vec{w}\vec{z}} = e^{i(kx+my+pz)}$$

$$\frac{\partial f}{\partial x} = e^{i\vec{w}\vec{z}} ik$$

$$\frac{\partial^2 f}{\partial x^2} = e^{i\vec{w}\vec{z}} i^2 k^2 = -e^{i\vec{w}\vec{z}} k^2$$

$$\begin{aligned} \therefore \nabla^2 f &= \left(\frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2} \right) \\ &= -e^{i\vec{w}\vec{z}} k^2 - e^{i\vec{w}\vec{z}} m^2 - e^{i\vec{w}\vec{z}} p^2 \\ &= e^{i\vec{w}\vec{z}} (-k^2 - m^2 - p^2) \\ &= -|\vec{w}|^2 e^{i\vec{w}\vec{z}} \end{aligned}$$

↓

$$\nabla^2 \phi(z) = -\frac{\rho(z)}{\epsilon_0}$$

$$\nabla^2 \int \tilde{\phi}(\vec{w}) e^{i\vec{w}\vec{z}} d\vec{z} = -\frac{1}{\epsilon_0} \int \tilde{\rho}(\vec{w}) e^{i\vec{w}\vec{z}} d\vec{z}$$

$$\int \tilde{\phi}(\vec{w}) \nabla^2 (e^{i\vec{w}\vec{z}}) d\vec{z} = -\frac{1}{\epsilon_0} \int \tilde{\rho}(\vec{w}) e^{i\vec{w}\vec{z}} d\vec{z}$$

$$\int \tilde{\phi}(\vec{w}) - |\vec{w}|^2 e^{i\vec{w}\vec{z}} d\vec{z} = -\frac{1}{\epsilon_0} \int \tilde{\rho}(\vec{w}) e^{i\vec{w}\vec{z}} d\vec{z}$$

$$|\vec{w}|^2 \tilde{\phi}(\vec{w}) = \frac{\tilde{\rho}(\vec{w})}{\epsilon_0}$$

附录[2]: 短程力推导

$$\text{令: } \vec{z} = [x, y, z]$$

$$f = \frac{\text{erfc}(\alpha |\vec{z} - \vec{z}_j + nL|)}{|\vec{z} - \vec{z}_j + nL|} = \frac{\text{erfc}(\alpha r)}{r}$$

$$r = |\vec{z} - \vec{z}_j + nL| = \sqrt{(x - x_j + nL)^2 + (y - y_j + nL)^2 + (z - z_j + nL)^2}$$

$$\text{则: } \frac{\partial r}{\partial x} = \frac{x - x_j + nL}{\sqrt{(x - x_j + nL)^2 + (y - y_j + nL)^2 + (z - z_j + nL)^2}} = \frac{x - x_j + nL}{r}$$

$$\begin{aligned} \frac{\partial f}{\partial x} &= \frac{\partial}{\partial r} \frac{\partial r}{\partial x} \left(\frac{1}{r} \right) \text{erfc}(\alpha r) + \frac{1}{r} \frac{\partial}{\partial r} \frac{\partial r}{\partial x} (\text{erfc}(\alpha r)) \\ &= \frac{\partial r}{\partial x} \left(\frac{\partial}{\partial r} \left(\frac{1}{r} \right) \text{erfc}(\alpha r) + \frac{1}{r} \frac{\partial}{\partial r} (\text{erfc}(\alpha r)) \right) \\ &= \frac{x - x_j + nL}{r} \left(-\frac{1}{r^2} \text{erfc}(\alpha r) - \frac{1}{r} \frac{2\alpha}{\sqrt{\pi}} e^{-\alpha^2 r^2} \right) \\ &= -\frac{x - x_j + nL}{r^3} \left(\text{erfc}(\alpha r) + \frac{2\alpha r}{\sqrt{\pi}} e^{-\alpha^2 r^2} \right) \end{aligned}$$

$$\begin{aligned} \nabla f &= \left(\frac{\partial f}{\partial x}, \frac{\partial f}{\partial y}, \frac{\partial f}{\partial z} \right) \\ &= - \left(\frac{x - x_j + nL}{r^3}, \frac{y - y_j + nL}{r^3}, \frac{z - z_j + nL}{r^3} \right) \left(\text{erfc}(\alpha r) + \frac{2\alpha r}{\sqrt{\pi}} e^{-\alpha^2 r^2} \right) \\ &= -\frac{\vec{z} - \vec{z}_j + nL}{r^3} \left(\text{erfc}(\alpha r) + \frac{2\alpha r}{\sqrt{\pi}} e^{-\alpha^2 r^2} \right) \end{aligned}$$

$$\therefore \nabla \phi_{[i]}^s = \sum_n \sum_{j=1(\neq i)}^N \nabla f q_j = - \sum_n \sum_{j=1(\neq i)}^N q_j \frac{\vec{z} - \vec{z}_j + nL}{r^3} \left(\text{erfc}(\alpha r) + \frac{2\alpha r}{\sqrt{\pi}} e^{-\alpha^2 r^2} \right)$$

附录[3]: 长程力推导

$$\vec{w} = [k, m, p], \vec{r} = [x, y, z]$$

$$\phi(r) = \Re \left(\sum_{\vec{w} \neq 0} \Phi(k) e^{i\vec{w}\vec{r}} \right)$$

$$\nabla_r \phi(r) = \nabla_r \Re \left(\sum_{\vec{w} \neq 0} \Phi(k) e^{i\vec{w}\vec{r}} \right)$$

$$= \Re \left(\nabla_r \sum_{\vec{w} \neq 0} \Phi(k) e^{i\vec{w}\vec{r}} \right)$$

$$= \Re \left(\nabla_{x,y,z} \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right)$$

$$= \left[\Re \left(\nabla_x \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right), \Re \left(\nabla_y \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right), \Re \left(\nabla_z \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right) \right]$$

$$= \left[\Re \left(ik \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right), \Re \left(im \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right), \Re \left(ip \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right) \right]$$

$$= \left[\Re \left(i \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right) k, \Re \left(i \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right) m, \Re \left(i \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right) p \right]$$

$$= \Re \left(i \sum_{\vec{w} \neq 0} \Phi(k) e^{i(kx+my+pz)} \right) [k, m, p]$$

$$= \Re \left(i \sum_{\vec{w} \neq 0} \Phi(k) e^{i\vec{w}\vec{r}} \right) \vec{w}$$

附录[4]

$$\begin{aligned}
& \int_{-\infty}^{\infty} e^{-\alpha^2 x^2} e^{-ikx} dx \\
&= \int_{-\infty}^{\infty} e^{-\alpha^2 x^2 - ikx + \frac{k^2}{4\alpha^2} - \frac{k^2}{4\alpha^2}} dx \\
&= e^{-\frac{k^2}{4\alpha^2}} \int_{-\infty}^{\infty} e^{-\alpha^2 x^2 - ikx + \frac{k^2}{4\alpha^2}} dx \\
&= e^{-\frac{k^2}{4\alpha^2}} \int_{-\infty}^{\infty} e^{-\left(\alpha x + \frac{ik}{2\alpha}\right)^2} dx \\
&= \frac{1}{\alpha} e^{-\frac{k^2}{4\alpha^2}} \int_{-\infty}^{\infty} e^{-\left(\alpha x + \frac{ik}{2\alpha}\right)^2} d\left(\alpha x + \frac{ik}{2\alpha}\right) \\
&= \frac{1}{\alpha} e^{-\frac{k^2}{4\alpha^2}} \int_{-\infty}^{\infty} e^{-y^2} dy \\
&= \frac{\sqrt{\pi}}{\alpha} e^{-\frac{k^2}{4\alpha^2}}
\end{aligned}$$

$$\begin{aligned}
& \int_{R^3} e^{-\alpha^2 |\vec{z} - \vec{z}_j|^2} e^{-i\vec{w}\vec{z}} d^3\vec{z} \\
&= e^{-i\vec{w}\vec{z}_j} \int_{R^3} e^{-\alpha^2 |\vec{z} - \vec{z}_j|^2} e^{-i\vec{w}(\vec{z} - \vec{z}_j)} d^3(\vec{z} - \vec{z}_j) \\
&= e^{-i\vec{w}\vec{z}_j} \int_{R^3} e^{-\alpha^2 |\vec{h}|^2} e^{-i\vec{w}\vec{h}} d^3\vec{h} \\
&= e^{-i\vec{w}\vec{z}_j} \iiint_{-\infty}^{\infty} e^{-\alpha^2(x^2+y^2+z^2)} e^{-i(kx+my+pz)} dx dy dz \\
&= e^{-i\vec{w}\vec{z}_j} \iiint_{-\infty}^{\infty} e^{-\alpha^2 x^2} e^{-\alpha^2 z^2} e^{-ikx} e^{-imy} e^{-ipz} dx dy dz \\
&= e^{-i\vec{w}\vec{z}_j} \int_{-\infty}^{\infty} e^{-\alpha^2 x^2} e^{-ikx} dx \int_{-\infty}^{\infty} e^{-\alpha^2 y^2} e^{-imy} dy \int_{-\infty}^{\infty} e^{-\alpha^2 z^2} e^{-ipz} dz \\
&= e^{-i\vec{w}\vec{z}_j} \frac{\pi^{\frac{3}{2}}}{\alpha^3} e^{-\frac{k^2+m^2+p^2}{4\alpha^2}} \\
&= e^{-i\vec{w}\vec{z}_j} \frac{\pi^{\frac{3}{2}}}{\alpha^3} e^{-\frac{|\vec{w}|^2}{4\alpha^2}}
\end{aligned}$$

$$\therefore \frac{\alpha^3}{\pi^{\frac{3}{2}}} \sum_{j=1}^N q_j \int_{R^3} e^{-\alpha^2 |\vec{z} - \vec{z}_j|^2} e^{-i\vec{w}\vec{z}} d^3\vec{z} = \frac{\alpha^3}{\pi^{\frac{3}{2}}} \sum_{j=1}^N q_j e^{-i\vec{w}\vec{z}_j} \frac{\pi^{\frac{3}{2}}}{\alpha^3} e^{-\frac{|\vec{w}|^2}{4\alpha^2}} = e^{-\frac{|\vec{w}|^2}{4\alpha^2}} \sum_{j=1}^N q_j e^{-i\vec{w}\vec{z}_j}$$

references

1. http://micro.stanford.edu/mediawiki/images/4/46/Ewald_notes.pdf
2. <http://staff.ustc.edu.cn/~zqj/posts/Ewald-Summation/#fn:FrenkelSmit>
3. <https://zhuatlan.zhihu.com/p/406664860>
4. <https://github.com/zhaodazhuang/fenxiang-/blob/main/Ewald%20Sum.pdf>